# Dynamic Pathways for Fluxional Molecules Defined Using Exchange-NOE Peaks 

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#### Abstract

In examining NMR methods to assess orientation and fluxional motions of planar N-donor heterocyclic coordinated ligands (L's) in solution, we introduce the use of exchange-NOE NMR data as a powerful method for defining (i) the extent of L rotation about the metal-N bond, (ii) the direction of L rotation, and (iii) even the halves of $C_{2}$-symmetrical L's that interchange during dynamic processes. The full value of the approach depends on a strategy in which the complexes studied are chiral and similar except that one has a lopsided $\mathrm{L}\left(\mathrm{Me}_{3} \mathrm{Bzm}=1,5,6-\right.$ trimethylbenzimidazole in $\left[\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}\right](\mathbf{1})$ ) and the other has a $C_{2}$-symmetrical $\mathrm{L}(3,5$-lut $=3,5$-lutidine in $\left[\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}(3,5-\mathrm{lut})_{4}\right](\mathbf{2})$. Each Re is part of a nearly linear $\mathrm{O}=\mathrm{Re}-\mathrm{O}-\mathrm{Re}=\mathrm{O}$ grouping and has a "terminal" L $\left(L^{t}\right)$ and a stacked $L\left(L^{s}\right)$. The fluxional inversion of the two chiral dimers involves rotations of $\sim 180^{\circ}$ about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds and of $\sim 90^{\circ}$ about all four $\mathrm{Re}-\mathrm{L}$ bonds. The exchange-NOE data for 2 show that the half of $\mathrm{L}^{\mathrm{t}}$ away from the dimer center interchanges with the half of $L^{s}$ close to the center, with the $L$ plane rotating past the $\mathrm{O}=\mathrm{Re}-\mathrm{O}$ bonds, not the $\mathrm{N}-\mathrm{Re}-\mathrm{Cl}$ bonds. Thus, the exchange-NOE data help to establish the direction of L rotation. In 1, $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ has the rare head-to-head (HH) orientation with respect to the partner. This partner $\mathrm{Me}_{3^{-}}$ $\mathrm{Bzm}^{\text {s }}$ stacks with the $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }}$ from the other Re in the common head-to-tail (HT) orientation. Compelling evidence that the predominant solution conformer of $\mathbf{1}$ has the $\mathrm{HH}, \mathrm{HT}, \mathrm{HH}$ structure includes unusual chemical shift dispersions and a strong interligand NOE cross-peak. This is the only case in which cis,bis imidazole-ring-ligated untethered ligands have been found to be predominantly HH in solution. This predominance can be attributed to the electrostatic attraction of the $\delta^{+} \mathrm{N}_{2} \mathrm{C}$ proton for the negative core of the molecule (bridging O , cis Cl on same Re , and two cis Cl 's on the other Re).


## Introduction

A wealth of solid-state structural information from X-ray crystallography has enriched our knowledge of the chemistry of metal compounds. In contrast, the dynamic pathways and the solution structures of fluxional species are much less well understood. However, it is usually in solution that metal species act as catalysts, as reagents for organic synthesis, as active sites in metalloenzymes, etc. The multitude of potential dynamic pathways that octahedral mononuclear and polynuclear metal complexes possess, combined with a paucity of suitable experimental approaches, has limited our knowledge of these dynamic processes. Here we present a general strategy that has provided excellent insight into the dynamic and structural properties of fluxional complexes in solution. Our study also reveals a feature of the coordinated imidazole ring that seems likely to influence the biological role of imidazole rings in many metallobiochemicals.

This study arose from an initial objective directed at understanding the factors influencing the orientations and dynamic motions of cis coordinated nucleopurines. These properties in metalloanticancer drug DNA adducts influence the nature of the adducts formed and the activity of the adducts. Nucleopurines, like other important N -donor heterocyclic bioligands, possess a plane of symmetry passing through the rings, but they are lopsided, i.e., they lack $C_{2}$-symmetry. Metalloproteins or enzymes in which the metal has two coordinated heterocyclic lopsided imidazole ligands are almost ubiquitous. ${ }^{1}$

Nucleobases have relatively few nonexchangeable protons, and their orientation in solution is difficult to define by NMR

[^0]spectroscopy. To circumvent this limitation, we have been using the nucleopurine analog, 1,5,6-trimethylbenzimidazole ( $\mathrm{Me}_{3}{ }^{-}$ Bzm, Chart 1). This ligand has the very useful characteristic of having protons distributed around the entire molecule; each proton has a singlet NMR signal. ${ }^{2}$ Like nucleopurines, $\mathrm{Me}_{3}{ }^{-}$ Bzm is a five-six bicyclic system but the six-membered ring lacks an exocyclic group directed toward the cis coordinated ligands. This difference in steric properties is helpful in assessing the contributions of steric effects to structural changes in DNA when metal species, such as metal anticancer agents, are bound to DNA. Finally, $\mathrm{Me}_{3} \mathrm{Bzm}$ is also of particular interest in $\mathrm{B}_{12}$ model chemistry. ${ }^{3}$ For two cis lopsided N heterocyclic ligands ( B, e.g., purines), the corresponding atoms of each B can be on the same or opposite sides of the $\mathrm{N}-\mathrm{M}-\mathrm{N}$ plane, giving the head-to-head (HH) and the head-to-tail (HT) orientations, respectively. In a preliminary communication, we described the rarity of finding the HH orientation of two cis lopsided ligands. ${ }^{2}$ A HT arrangement is the most common solidstate conformation of cis-bis(ligand) complexes of purines with $\mathrm{Pt}^{\mathrm{II}}, \mathrm{Co}^{\text {III }}, \mathrm{Cu}^{\text {II }}$, and $\mathrm{Zn}^{\text {II }}$ metal centers. ${ }^{4}$ Only four HH solidstate structures of untethered ligands had been reported, all of which are square-planar $\mathrm{Pt}^{\mathrm{II}}$ complexes with 9 -ethylguanine, ${ }^{5}$ until our recent brief report on the dynamic properties and stereochemistry of $\mathrm{Ru}-$ and $\mathrm{Re}-\mathrm{Me}_{3} \mathrm{Bzm}$ complexes. ${ }^{2}$ For one of these rare HH species, $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}(\mathbf{1})$, only a figure of the structure was shown. ${ }^{2}$ Our analysis of new results on this and related Re species has led to the identification of some interesting and important new phenomena which are likely to

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Chart 1. NMR Labeling Scheme for the $\mathrm{Me}_{3} \mathrm{Bzm}$ Ligand $^{a}$

${ }^{a}$ The carbons are designated as " B " which refers to the benzimidazole ring.

Table 1. Crystallographic Data for
$\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4} \cdot 2\left(\mathrm{CH}_{3}\right)_{2} \mathrm{CO}$

| chemical formula | $\mathrm{C}_{46} \mathrm{H}_{60} \mathrm{Cl}_{4} \mathrm{~N}_{8} \mathrm{O}_{5} \mathrm{Re}_{2}$ |
| :--- | :--- |
| formula weight | 1319.2 |
| space group | $C 2 / c$ |
| $Z$ | 8 |
| $a(\AA)$ | $35.020(12)$ |
| $b(\AA)$ | $16.308(5)$ |
| $c(\AA)$ | $25.507(7)$ |
| $\beta(\mathrm{d} \mathrm{\AA})$ | $132.870(0)$ |
| $V\left(\AA^{3}\right)$ | $10676(8)$ |
| $\rho_{\text {caldd }}\left(\mathrm{g} / \mathrm{cm}^{3}\right)$ | 1.64 |
| $\lambda(\AA)$ | $0.71073(\mathrm{Mo} \mathrm{K} \mathrm{\alpha})$ |
| $\mu\left(\mathrm{mm}^{-1}\right)$ | 4.84 |
| $T(\mathrm{~K})$ | 173 |
| scan technique | $\omega$ |
| transm coeff $(\%)$ | $4.7-11.9$ |
| $R(\%)$ | 5.62 |
| $R_{w}(\%)$ | 6.43 |
| goodness-of-fit | 1.43 |

be significant in areas beyond those which motivated our study. Furthermore, our study shows the power of the interpretation of exchange-NOEs in providing insight into the dynamic motions of molecules.

## Experimental Section

$\mathbf{R e}_{2} \mathbf{O}_{3} \mathbf{C l}_{\mathbf{4}}\left(\mathbf{M e}_{3} \mathbf{B z m}\right)_{\mathbf{4}} \mathbf{( 1 )}$. A solution of $\mathrm{Me}_{3} \mathrm{Bzm}(0.26 \mathrm{~g}, 1.6 \mathrm{mmol})^{3}$ in acetone ( 20 mL containing $10 \mu \mathrm{~L}$ of $\mathrm{H}_{2} \mathrm{O}$ ) was added dropwise to a stirred suspension of $\mathrm{ReOCl}_{3}\left(\mathrm{Me}_{2} \mathrm{~S}\right)\left(\mathrm{OPPh}_{3}\right)^{6}(0.325 \mathrm{~g}, 0.5 \mathrm{mmol})$ in acetone $(30 \mathrm{~mL})$. The mixture was left at room temperature overnight. The emerald-green solution obtained after filtration of a white precipitate $\left(\mathrm{Me}_{3} \mathrm{Bzm} \cdot \mathrm{HCl}\right)$ was concentrated to $\sim 15 \mathrm{~mL}$ and refrigerated. A deepgreen crystalline precipitate, formed within 2 days, was collected, washed with water, cold acetone, and diethyl ether, and then vacuum dried: yield $0.18 \mathrm{~g}(60 \%)$. Anal. Calcd for $\mathrm{C}_{40} \mathrm{H}_{48} \mathrm{Cl}_{4} \mathrm{~N}_{8} \mathrm{O}_{3} \mathrm{Re}_{2}\left(M_{\mathrm{w}}\right.$ 1203.09): C, 39.93; H, 4.02; N, 9.31; Cl, 11.78. Found: C, 39.90; H, $4.06 ; \mathrm{N}, 9.22 ; \mathrm{Cl}, 11.86$. The complex $(0.12 \mathrm{~g})$ was recrystallized from dimethylformamide (DMF)/acetone mixtures $(0.9 \mathrm{~mL} / 6 \mathrm{~mL})$ by addition of diethyl ether. The crystals contain acetone (by NMR) but desolvate readily and could not be analyzed: yield $65 \%$.
$\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathbf{3}, 5\right.$-lut) $\mathbf{4}_{\mathbf{~ ( 2 )}}$ ( (lut $=$ lutidine). 2 was prepared by the method used for $\mathbf{1}$. The clear emerald-green solution was concentrated to $\sim 10 \mathrm{~mL}$ and then refrigerated. A deep-green precipitate formed in a few hours and was collected in $72 \%$ yield and dried as above. Anal. Calcd for $\mathrm{C}_{28} \mathrm{H}_{36} \mathrm{Cl}_{4} \mathrm{~N}_{4} \mathrm{O}_{3} \mathrm{Re}_{2}\left(M_{\mathrm{w}} 990.84\right)$ : C, 33.94; H, 3.66; $\mathrm{N}, 5.65$; Cl, 14.31. Found: C, 34.16 ; H, 3.78; N, 5.55; Cl, 14.21.

X-ray Crystallography. Crystal data and refinement parameters are summarized in Table 1. A green crystal of $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3}-\right.$ $\mathrm{Bzm})_{4} \cdot 2 \mathrm{CH}_{3} \mathrm{COCH}_{3}(\mathbf{1})$ was mounted on a glass fiber, placed on a Siemens P4 diffractometer, and cooled to $-100^{\circ} \mathrm{C}$ with a LT2 lowtemperature device to prevent crystal decay. Crystals of $\mathbf{1}$ desolvate within 20 min when removed from the mother liquor at room temperature. The crystal system for $\mathbf{1}$ was determined by automatic indexing of 38 centered reflections. Refined cell constants were obtained by least-squares refinement of 15 of the high-angle reflections $\left(17^{\circ}<2 \theta<44^{\circ}\right)$. Intensities were collected using monochromated
(6) Grove, D. E.; Wilkinson, G. J. Chem. Soc. A 1966, 1224-30.


Figure 1. Perspective drawing of $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}$ (1) with $50 \%$ probability for the thermal ellipsoids.

Mo $\mathrm{K} \alpha$ radiation $(\lambda=0.71073 \AA)$. Three check reflections were measured every 97 reflections; there was no significant deviation in intensity. Intensities were corrected for Lorentz and monochromator polarization effects and absorption (semiempirical method based on azimuthal scans). Systematic absences were consistent with space groups $C 2 / c$ (centrosymmetric) and $C c$ (noncentrosymmetric). The space group $C 2 / c$ was selected on the basis of the statistical distribution of intensities (mean $\left|E^{2}-1\right|=1.030$ ). The structure was solved by Patterson methods and refined by full-matrix least-squares procedures using SHELXTL PLUS (VMS). The acetone molecules in $\mathbf{1}$ were highly disordered and were refined isotropically as rigid groups pivoting about the central carbon atoms ( $\mathrm{C} 41, \mathrm{C} 44$ ) with idealized bond lengths $(d(\mathrm{C}-\mathrm{O})=1.302 \AA, d(\mathrm{C}-\mathrm{C})=1.540 \AA)$ and bond angles $\left(120^{\circ}\right)$. All other non-hydrogen atoms were refined anisotropically. H -atoms were included at calculated $(d(\mathrm{C}-\mathrm{H})=0.96 \AA)$ positions with isotropic thermal parameters fixed at 0.08 .

NMR Spectroscopy. NMR spectrometers used were JEOL EX400 and General Electric QE-300, NT-360, and GN-600 Omega models. Sample concentrations were $\sim 50 \mathrm{mM}$. All spectra were referenced to TMS unless otherwise noted. ${ }^{1} \mathrm{H}$ 1D data were collected typically with a $3600-7200 \mathrm{~Hz}$ window, a $30^{\circ}$ pulse width, and 16 K data points. Variable temperature (VT) 2D NOESY/EXSY spectra were recorded on the JEOL instrument. 2D parameters include: matrix of $512 \times$ 1024; the second dimension was zero filled to 1024 data points prior to Fourier transformation; eight scans per $t_{1}$ increment were preceded by two dummy scans, with a spectral window 3959 Hz , pulse delay of 2.37 s , a mixing time of 500 ms , and a square sine bell function with no phase shift applied in both dimensions.

## Results

Structural Studies. A perspective drawing of the $\mathrm{Re}_{2} \mathrm{O}_{3}-$ $\mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}(\mathbf{1})$ structure appears in Figure 1. As in the py (pyridine) analogue, $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}(\mathrm{py})_{4}$ (3), ${ }^{7}$ the nitrogen donor ligands on each pseudo-octahedral Re atom are cis to each other, and the molecule has a pseudo two-fold axis through the bridging oxo ligand, $\mathrm{O}(2)$. The most interesting features of the molecule are the HH arrangement of the two cis $\mathrm{Me}_{3} \mathrm{Bzm}$ ligands on each Re (the "Re-ligand-partners") and the extended network of four $\mathrm{Me}_{3} \mathrm{Bzm}$ ligands. The center of the network has two HT pseudo-bridging "stacked" Me ${ }_{3}$ Bzm's, designated as $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$. The Re-ligand-partner of each $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ is designated $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ for "terminal".

In discussing the $\mathrm{Me}_{3} \mathrm{Bzm}$ parameters, we use the numbers in Chart 1 and refer to carbons as B2, etc. All four $\mathrm{Re}-\mathrm{N}-\mathrm{B} 2$ bond angles (supporting information) are consistently more acute than the four $\mathrm{Re}-\mathrm{N}-\mathrm{B} 9$ bond angles. Unsymmetrical angles are common in $\mathrm{Co}^{\text {III }}$ complexes with $\mathrm{Me}_{3} \mathrm{Bzm}$ : $\mathrm{Co}-\mathrm{N}-\mathrm{B} 2$


Figure 2. Perspective view of the intramolecular stacking interaction between the $\mathrm{Me}_{3} \mathrm{Bzm}^{5}$ ligands for $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}$ (1). The two $\mathrm{Me}_{3} \mathrm{Bzm}^{t}$ ligands are omitted for clarity.
angles range $121.1(3)-123.1(2)^{\circ} ; \mathrm{Co}-\mathrm{N}-\mathrm{B} 9$ angles range 132.3(3)-133.7(2) ${ }^{\circ} .{ }^{3,8-10}$

The $\mathrm{Me}_{3} \mathrm{Bzm}$ ligands in $\mathbf{1}$ are rotated about the $\mathrm{Re}-\mathrm{N}$ bond $34-54^{\circ}$ away from coplanarity with the plane defined by Re and the cis Cl atoms. In addition, the two halves of the dimer are rotated about the $\operatorname{Re}(1)-\operatorname{Re}(2)$ vector away from an eclipsed conformation. Treating the $\mathrm{Re}-\mathrm{Re}$ vector as a "bond", we find that a $\mathrm{N}(3)-\operatorname{Re}(1)-\operatorname{Re}(2)-\mathrm{N}(7)$ "torsion angle" of $28^{\circ}$ results from the greater overlap of the $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ ligands (Figure 2). The shortest interplanar distance between the $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ ligands is $3.1 \AA$. The two $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ ligands are rotated $\sim 180^{\circ}$ away from each other.

NMR Spectroscopy. The 1D and 2D ${ }^{1} \mathrm{H}$ NMR spectra of $\mathbf{1}$ and 2 (Figures 3-6) are consistent with both being fluxional molecules. At $20^{\circ} \mathrm{C}, \mathbf{1}$ has only six signals for the $\mathrm{Me}_{3} \mathrm{Bzm}$ ligand, and these signals are somewhat broad, suggesting an exchange process with a rate at the fast end of the intermediate NMR time scale. For 2, the signals are sharp, consistent with a faster exchange process. Spectra at lower temperature for $\mathbf{1}$ and 2 (Figures 3-6) show initial broadening of peaks. At very low temperatures, the spectra for both sharpen with the emergence of nearly twice the number of peaks apparent at room temperature. These results suggest that the molecules exist in solution with the N -heterocycle in two different environments and that the ligands interchange position between these environments relatively rapidly at ambient temperatures. Details of the NMR results are discussed below.

## Discussion

Solid-State Structure. The interesting extended HH,HT,HH arrangement of $\mathrm{Me}_{3} \mathrm{Bzm}$ 's in 1 (Figure 1) can be attributed in large part to the electrostatic attraction of the $\delta^{+} \mathrm{H} 2$ 's on $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ 's toward the negative O and Cl groups in the core of the dimer. There are several reasons why we believe this is the key interaction that dictates conformation. First, there are so many favorable interactions; for example, the $\delta^{+} \mathrm{H}$ on $\mathrm{C}(21)$ in Figure 1 is attracted by four negative groups, $\mathrm{Cl}(4), \mathrm{Cl}(2)$, $\mathrm{Cl}(1)$, and $\mathrm{O}(2)$. Second, in a related dimer with one pyridine (py) and one $\mathrm{Me}_{3} \mathrm{Bzm}$ on each Re, the py occupies the pseudo bridging site while the $\mathrm{Me}_{3} \mathrm{Bzm}$ occupies the terminal site with the same orientation found in $\mathbf{1} .{ }^{11}$ Py lacks the very acidic H characteristic of the H 2 in an imidazole ring. Finally, although
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Figure 3. ${ }^{1} \mathrm{H}$ NMR spectra of $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}$ (1) in $\mathrm{CD}_{2} \mathrm{Cl}_{2}$ at various temperatures in the downfield $(\mathrm{A})$ and upfield ( B ) regions. Signal assignments for the stacked and terminal $\mathrm{Me}_{3} \mathrm{Bzm}$ are indicated.


Figure 4. $2 \mathrm{D}^{1} \mathrm{H}$ NOESY/EXSY spectrum of 1 in $\mathrm{CD}_{2} \mathrm{Cl}_{2}$ at $-90^{\circ} \mathrm{C}$. The exchange pathways are shown. Some exchange-NOE peaks are marked with an asterisk, for example, the $2 \mathrm{t}-12 \mathrm{~s}$ exchange-NOE.
$\mathrm{Re}^{\mathrm{V}}$ and $\mathrm{Ru}^{\mathrm{II}}$ are very different metal centers, the relative stability of the atropisomers found for $\mathrm{Ru}^{\mathrm{II}}$ complexes of lopsided ligands can best be rationalized by invoking this type of electrostatic attraction of the $\delta^{+} \mathrm{H} 2$ toward the coordinated Cl ligands. ${ }^{12}$
The distribution of negative groups surrounding the stacked site is more symmetrical than for the terminal site. In the observed $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ orientation, the bulky $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ sixmembered ring is directed away from the bulkiest ligand on the same $\mathrm{Re}, \mathrm{Me}_{3} \mathrm{Bzm}^{t}$. The lower steric crowding favors the observed $\mathrm{HT} \mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }}$ orientation compared to the orientation obtained by rotation of $\mathrm{Me}_{3} \mathrm{Bzm}^{5}$ by $180^{\circ}$ from that in 1. These steric and electrostatic factors act in concert to favor the $\mathrm{HH},-$ HT,HH conformer observed in the solid.

Spectral Assignments and Solution State Structure. The spectral assignments and the assessment of the structure of $\mathbf{1}$ in $\mathrm{CD}_{2} \mathrm{Cl}_{2}$ depended on anisotropic shift arguments and on EXSY cross-peaks in the 2D NOESY/EXSY data at $-90^{\circ} \mathrm{C}$ (Figure 4). As background for our assignment, we first review our use of the anisotropic effects of one $\mathrm{Me}_{3} \mathrm{Bzm}$ on the shifts of the signals of the other $\mathrm{Me}_{3} \mathrm{Bzm}$ to assign signals and determine the relative $\mathrm{Me}_{3} \mathrm{Bzm}$ orientations of cis,cis,cis$\mathrm{RuCl}_{2}\left(\mathrm{Me}_{2} \mathrm{SO}_{2}\right)_{2}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{2} .^{2}$ The anisotropic effects were greatest on the H 2 and H 4 signals. In the canted orientation adopted by HH species, H 2 of one $\mathrm{Me}_{3} \mathrm{Bzm}$ is pointed toward the other $\mathrm{Me}_{3} \mathrm{Bzm}$ and H 4 of this other $\mathrm{Me}_{3} \mathrm{Bzm}$ is pointed toward the first $\mathrm{Me}_{3} \mathrm{Bzm}$. As a result, a characteristic up/ downfield set of H 2 and of H 4 signals is produced. Since H7 of both $\mathrm{Me}_{3} \mathrm{Bzm}$ points away from the coordination sphere, the anisotropic effects on this signal were too small to aid structure elucidation; the H 7 assignments depended on NOE data. The NOE patterns readily allowed signal assignments since the $\mathrm{Me}_{3}{ }^{-}$ Bzm ligands occupy distinct sites in the Ru ${ }^{\text {II }}$ complex and do not interchange position.
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$-60{ }^{\circ} \mathrm{C}$
$-70{ }^{\circ} \mathrm{C}$
$-80^{\circ} \mathrm{C}$


Figure 5. ${ }^{1} \mathrm{H}$ NMR spectra of $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}(3,5-\mathrm{lut})_{4}$ (2) in $\mathrm{CD}_{2} \mathrm{Cl}_{2}$ at various temperatures in the downfield (A) and upfield (B) regions. Signal assignments for the stacked and terminal 3,5-lut are indicated.

The $\mathrm{Me}_{3} \mathrm{Bzm}$ interchange between positions in $\mathbf{1}$ makes the 2 D assignment procedure more difficult than for the Ru compound. We made an initial assumption that the conforma-


Figure 6. 2D ${ }^{1} \mathrm{H}$ NOESY/EXSY spectrum of 2 in $\mathrm{CD}_{2} \mathrm{Cl}_{2}$ at -100 ${ }^{\circ} \mathrm{C}$. Since the combined rotations around $\mathrm{Re}-\mathrm{O}$ and $\mathrm{Re}-\mathrm{N}$ bonds will exchange all signals of a given type, for each $\alpha-\mathrm{H}$ and $\beta-\mathrm{CH}_{3}$ signal, there are three exchange cross-peaks to the other three $\alpha-\mathrm{H}$ and $\beta-\mathrm{CH}_{3}$ signals, respectively. The $\gamma$-H signals exchange also. These cross-peaks are close to the diagonal. Cross-peaks far from the diagonal are between $\beta-\mathrm{CH}_{3}$ and $\alpha-/ \gamma-\mathrm{H}$ signals and are shown at two levels. The lower left of the figure (at a higher cut) shows only direct NOEs from $\beta-\mathrm{CH}_{3}$ to $\alpha-\mathrm{H}$ or $\gamma-\mathrm{H}$ signals. Higher cuts allowed assignments and allowed us to identify the direct NOEs. In the upper right (an intermediate contour), there are exchange-NOE peaks from $\beta-\mathrm{CH}_{3}$ to $\gamma-\mathrm{H}$ (class $\mathrm{i}, *$ ) and $\alpha-\mathrm{H}$ (class ii, ${ }^{* *}$ ) signals marked with asterisks. However, there are two other exchange-NOEs (class iii, ****). These are from the upfield $_{\text {the }}$ $\alpha-\mathrm{H}$ signal of the $3,5-\mathrm{lut}$ to the downfield $\beta-\mathrm{CH}_{3}$ signal of $3,5-\mathrm{lut}{ }^{t}$ and from the downfield $\alpha-\mathrm{H}$ signal of $3,5-1 \mathrm{lut}^{\mathrm{t}}$ to the upfield $\beta-\mathrm{CH}_{3}$ signal of $3,5-\mathrm{lut}$.
tion of $\mathbf{1}$ in the crystal also exists in solution at low temperature, except of course that $\mathbf{1}$ will have true $C_{2}$ symmetry and thus is chiral. We considered alternative structures, but only our initial assumption was totally consistent with all the NMR data.

In the assignments at $-90^{\circ} \mathrm{C}$ (Table 2), signals are labeled by the position of the proton and the type of $\mathrm{Me}_{3} \mathrm{Bzm}$. The first step in the assignment was to link protons of a given type in the two $\mathrm{Me}_{3} \mathrm{Bzm}$ 's using the EXSY peaks, which are very clear and the only peaks present at the highest contours that contain cross-peaks. The pattern, two pairs of one upfield and one downfield exchanging aromatic signals, is characteristic of B 4 H and B 2 H signals of an HH species. ${ }^{2}$ No HT species is known to give this pattern. We assign the most downfield signal as 2 t , the signal of B 2 H in $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$. B 2 H is ordinarily the most downfield $\mathrm{Me}_{3} \mathrm{Bzm}$ signal, and B 2 H in $\mathrm{Me}_{3} \mathrm{Bzm}^{t}$ is far from the shielding region of the anisotropic $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$. 2t has an EXSY peak to the upfield 2 s , thus assigning it. In the other well-dispersed EXSY pair, the far upfield shifted signal must be 4 t since B 4 H of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ is directed toward the shielding region of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$.

As expected on the basis of the HH structure, the most downfield methyl signal is assigned to 12 t , since the $\mathrm{B} 12 \mathrm{H}_{3}$ group of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ is attached to a nitrogen and is directed away from the shielding region of the $\mathrm{Me}_{3} \mathrm{Bzm}^{5}$. The upfield EXSY
partner $\left(\mathrm{B} 12 \mathrm{H}_{3}\right.$ on $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }}$ directed toward the shielding region of $\left.\mathrm{Me}_{3} \mathrm{Bzm}^{t}\right)$ is usually found at $\sim 3.3 \mathrm{ppm}(\sim 0.5 \mathrm{ppm}$ from the typical shift of a coordinated $\mathrm{Me}_{3} \mathrm{Bzm}$ ). ${ }^{2}$ However, the 12s signal is at $2.37 \mathrm{ppm}, \sim 1 \mathrm{ppm}$ even further upfield. This unusual additional upfield shift can be explained readily with the $\mathrm{HH}, \mathrm{HT}, \mathrm{HH}$ structure. Each $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}} \mathrm{B} 12 \mathrm{H}_{3}$ group lies over and is shielded by the six-membered ring of the $C_{2}{ }^{-}$ related $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ on the other Re of the dimer (Figure 2). The most upfield signal is $10 t$ since the $B 10 H_{3}$ group of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ lies somewhat over $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }}$. Thus, the NMR shift data are totally consistent with the same structure for $\mathbf{1}$ in the solid and solution states, except that a more symmetrical version exists in solution.

Next we used the 2D data to confirm the assignments based on shift arguments and to assign the signals with relatively small dispersion between the two types of $\mathrm{Me}_{3} \mathrm{Bzm}$. The NOESY/ EXSY spectra of $\mathbf{1}$ (see Figure 4) have three types of crosspeaks. The first type, EXSY cross-peaks found at high contours and generally in the vicinity of the diagonal, was mentioned above as resulting from the physical exchange of the same $\mathrm{Me}_{3}-$ Bzm between the stacked and the terminal sites. The other two types are smaller and generally far from the diagonal. The second type is a direct NOE, e.g., 4t-10t. The third type, an exchange-NOE, results from direct NOE interactions followed by physical exchange of the ligands between the two types of sites, e.g., $4 \mathrm{~s}-10 \mathrm{t}$. This exchange-NOE information is useful for gaining a deeper understanding of the pathways for exchange in 2, see below.

The 2D data for $\mathbf{1}$ are somewhat complicated. First, the $t-t$ NOE cross-peaks are generally stronger than the s-s cross-peaks. This may be due to differences in correlation times for the different parts of the molecule. Second, the 11s and 10s signals overlap. The complete assignment required analysis of peak intensities at several mixing times ( 150,300 , and 500 ms ) and depended primarily on the better resolved t signals. At the 150 ms mixing time, the direct NOEs are most clearly identified. There are clear $t-t$ NOE cross-peaks linking adjacent groups, $4 \mathrm{t}-10 \mathrm{t}-11 \mathrm{t}-7 \mathrm{t}-12 \mathrm{t}-2 \mathrm{t}$. The NOE data allow confirmation of the assignments from shift arguments and the assignment of the two less dispersed signals, 11t and 7 t . The t -s EXSY peaks allowed assignment of the s signals.

All the $t$-t NOE cross-peaks, except for $7 \mathrm{t}-12 \mathrm{t}$, appear before any s-s NOE cross-peaks. As can be seen from Chart 1, the intraligand H 7 to $\mathrm{B} 12 \mathrm{H}_{3}$ distance is relatively large since these protons reside on different rings. Thus the intraligand 7 to 12 NOE should be smaller than the other NOEs, accounting for the weak $7 \mathrm{t}-12 \mathrm{t}$ NOE. However, the $7 \mathrm{~s}-12 \mathrm{~s}$ NOE is one of the two strongest s-s cross-peaks. In the HT stacked arrangement, each $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }} \mathrm{H} 7$ is close to two $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }} \mathrm{B} 12 \mathrm{H}_{3}$ groups. Thus, the resulting two pathways for spin-spin interactions lead to a relatively large $7 \mathrm{~s}-12 \mathrm{~s}$ NOE cross-peak. The large $7 \mathrm{~s}-12 \mathrm{~s}$ NOE cross-peak strongly supports the conclusion from shift arguments that, in solution, $\mathbf{1}$ has the $\mathrm{HH}, \mathrm{HT}, \mathrm{HH}$ structure.

Since the exchange-NOE peaks result from a build-up of a direct-NOE followed by a dynamic process, they usually are seen most easily at longer mixing times. However, since the direct-NOEs also decay over time, and since there are many factors influencing the build-up and decay rates, it is not straightforward to treat the data quantitatively. At 500 ms , some exchange-NOE peaks can be observed at higher contour levels than the related direct-NOE peak. The relatively large 7s-12s NOE cross-peak leads to a $7 \mathrm{~s}-12 \mathrm{t}$ exchange-NOE cross-peak that is larger than the direct 7t-12t NOE peak (Figure 4). Of the first eight non-EXSY peaks that we observe at the higher contours, one is an exchange-NOE at 150 ms but four are exchange-NOEs at 500 ms .



Figure 7. Schematic representation of the pairwise relationship of the L ligands during $180^{\circ}$ net rotation about one $\mathrm{Re}-\mathrm{O}$ bond $\left(\mathrm{Re}_{\mathrm{b}}-\mathrm{O}\right)$ in $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4} \mathrm{~L}_{4}$ complexes. View is along the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds with $\mathrm{Re}_{\mathrm{a}}$ obscuring $\mathrm{Re}_{\mathrm{b}}$. The pairs, (1a and 1b) and (2a and 2b), interchange between occupying both terminal sites and both stacked sites.

## Dynamic Processes in $\left[\mathrm{ReOCl}_{2} \mathbf{L}^{s} \mathbf{L}^{\dagger}\right] \mathrm{O}\left[\mathrm{ReOCl}_{2} \mathrm{~L}^{s} \mathrm{~L}^{\dagger}\right]$ Dimers.

 As mentioned, the ${ }^{1} \mathrm{H}$ NMR spectra of $\mathbf{1}$ and $\mathbf{2}$ at $20^{\circ} \mathrm{C}$ indicate that both are dynamic. It is worthwhile to consider the general types of processes which can occur. The NMR data for 2, discussed last, are particularly revealing and illustrate the power of using a combination of exchange-NOE information and $C_{2}-$ symmetrical ligands to elucidate dynamic processes.The types of dynamic processes that could be important are (a) rotation about the $\mathrm{Re}-\mathrm{N}$ bonds and (b) rotation about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds. The latter process is shown schematically in Figure 7. Combinations of (a) and (b) can occur. In order for the $L^{\mathrm{s}}$ and $\mathrm{L}^{\mathrm{t}}$ signals to average, there must be rotation about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds. This process must involve either two $\sim 90^{\circ}$ rotations (one about each $\mathrm{Re}-\mathrm{O}$ bond) or one $\sim 180^{\circ}$ rotation about one $\mathrm{Re}-\mathrm{O}$ bond. The $90^{\circ}$ rotations could occur either stepwise or simultaneously. In either case, the L ligands can be seen to act in a pairwise fashion, unless there is bond breaking. Pair 1 (one L pair-mate on each Re ) can together interchange position, both being either the $L^{s}$ or the $L^{t}$ type (Figure 7). Pair 2 (again one L pair-mate on each Re) can interchange likewise, always occupying the opposite type of position to pair 1. The pairwise relationship just described holds true even if the symmetry of an intermediate were to make the four ligands equivalent. The magnetic equivalence of protons does not mean they physically interchange; it means only that they interchange positions.
The overall dynamic process inverts the chirality of the dimer. The simple rotation of a net $180^{\circ}$ around the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds converts the groupings $\mathrm{L}^{\mathrm{s}}-\mathrm{Re}-\mathrm{Cl}$ and $\mathrm{L}^{\mathrm{t}}-\mathrm{Re}-\mathrm{Cl}$ to $\mathrm{L}^{\mathrm{t}}-\mathrm{Re}-$ Cl and $\mathrm{L}^{\mathrm{s}}-\mathrm{Re}-\mathrm{Cl}$, respectively. It can be seen that the dimers in Figure 7 are mirror images of each other. The L ligand planes lie between the adjacent bonds, not along them. Such orientations are a component of the canting of one L toward the other. In 1, H 2 of $\mathrm{Me}_{3} \mathrm{Bzm}^{t}$ lies along the direction of the plane approximately bisecting the bonds to the cis Cl and the cis bridging O , but H 2 of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{s}}$ has a similar relationship to the bonds to the cis $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ and the cis bridging O . Thus, at some point during the inversion, each of the four $\mathrm{Me}_{3} \mathrm{Bzm}$ 's must also rotate by a net $\sim 90^{\circ}$ about the $\mathrm{Re}-\mathrm{N}$ bond (Figure 8).

Four types of $\mathrm{Re}-\mathrm{N}$ rotations of $180^{\circ}$ or less can be envisioned: clockwise by $\sim 90^{\circ}$ and $\sim 180^{\circ}$ or counterclockwise by $\sim 90^{\circ}$ and $\sim 180^{\circ}$. We will show that the $\mathrm{Re}-\mathrm{N}$ rotation process can be assessed with unprecedented insight by using the combined results from complexes with $C_{2}$-symmetrical ligands and from complexes with lopsided ligands.

Dynamic Processes in 1. The averaging of each type of $\mathrm{Me}_{3} \mathrm{Bzm}$ signal for the two types of $\mathrm{Me}_{3} \mathrm{Bzm}$ ligands of $\mathbf{1}$ at $20{ }^{\circ} \mathrm{C}$ (Figure 3) requires that rotation about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds is rapid. However, the $\mathrm{Me}_{3} \mathrm{Bzm}$ ligands could either rotate freely about the $\mathrm{Re}-\mathrm{N}$ bonds or maintain a fixed orientation (i.e., the Re-ligand-partners remain HH with each rotating by $\sim 90^{\circ}$ during the inversion).


Figure 8. Schematic representation showing the relationship of the plane of a lopsided ligand (represented by the arrow) to the different CP and comparing the simple $90^{\circ}$ rotation pathway to the more complicated $270^{\circ}$ rotation pathway.

At lower temperature, broadening was observed for all resonances, including the sharper B 7 H and $\mathrm{B} 11 \mathrm{H}_{3}$ resonances (Figure 3). At $-40{ }^{\circ} \mathrm{C}$ the $\mathrm{B} 2 \mathrm{H}, \mathrm{B} 4 \mathrm{H}, \mathrm{B} 12 \mathrm{H}_{3}$, and $\mathrm{B} 10 \mathrm{H}_{3}$ signals collapsed nearly into the baseline. As mentioned above, at lower temperature each signal split into two signals. Although slight changes in shift with temperature are common, the average shift of the two signals arising from each $20^{\circ} \mathrm{C}$ signal was close to that of the unsplit signal. Since other conformations would not have the same distinctive pattern of shifts, the similarity of the shifts at low and high temperatures suggests that the major conformer of $\mathbf{1}$ has the $\mathrm{HH}, \mathrm{HT}, \mathrm{HH}$ structure regardless of temperature. Any intermediate species formed in the rotation process is likely to be unstable with respect to the ground state, since otherwise its signals would average with those of the major conformer and the observed shifts would be temperature dependent.

The EXSY peaks in the 2D spectrum (Figure 4) could arise solely from the inversion of chirality. Indeed, in contrast to the results with the 3,5 -lut dimer ( $\mathbf{2}$, see below), the signals for the $\mathrm{Me}_{3} \mathrm{Bzm}$ dimer are equally broadened at -60 to $-70^{\circ} \mathrm{C}$. In the limit of slow exchange, signal broadening depends only on the rate constant for the process and not on the frequency difference between the exchanging signals. ${ }^{13}$ The essentially equal broadening of the signals of both $\mathrm{Me}_{3} \mathrm{Bzm}$ 's indicates that the broadening of both sets is due to one rate-limiting exchange process.

In summary, all the NMR observations for $\mathbf{1}$ are consistent with (a) the presence of only one major conformer, (b) a HH ,HT,HH structure for the major conformer, and (c) only one ratelimiting process.

Dynamic Processes in 2. For a $C_{2}$-symmetrical ligand, the exchange-NOE signals provide a high level of insight into dynamic processes. Furthermore, unlike the situation for lopsided ligands, the number of NMR signals expected depends on the extent of rotation about the $\mathrm{M}-\mathrm{N}$ bonds $\left(90^{\circ}\right.$ or $\left.180^{\circ}\right)$ and the rate of this process.

If rotation about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bonds and the net $90^{\circ} \mathrm{Re}-\mathrm{N}$ rotation accompanying inversion of chirality are rapid in $\mathrm{Re}_{2} \mathrm{O}_{3}{ }^{-}$ $\mathrm{Cl}_{4}(3,5-\mathrm{lut})_{4}$ (2), the signals expected for the $C_{2}$-symmetrical 3,5-lut ligand depend on the rate of the $180^{\circ}$ rotation about the $\mathrm{Re}-\mathrm{N}$ bonds. For very slow $180^{\circ}$ rotation, 3,5 -lut would give
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Table 2. Low-temperature ${ }^{1} \mathrm{H}$ NMR Assignments and Shifts for $\mathbf{1}$ and $\mathbf{2}^{a}$

| $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}\left(\mathrm{Me}_{3} \mathrm{Bzm}\right)_{4}(\mathbf{1})$ at $-90{ }^{\circ} \mathrm{C}$ |  |  | $\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4}(3,5-\mathrm{lut})_{4}(\mathbf{2})$ at $-100{ }^{\circ} \mathrm{C}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {t }}$ | $\mathrm{Me}_{3} \mathrm{Bzm}^{\text {s }}$ |  | 3,5-lut ${ }^{\text {t }}$ | 3,5-lut ${ }^{\text {s }}$ |
| B2H | 8.83 | 7.25 | $\alpha-\mathrm{H}^{\text {b }}$ | 7.38 | 7.22 |
| B4H | 6.29 | 7.96 | $\beta-\mathrm{CH}_{3}{ }^{\text {b }}$ | 2.05 | 1.73 |
| B7H | 7.28 | 7.05 | $\gamma$-H | 7.43 | 7.08 |
| $\mathrm{B} 10 \mathrm{H}_{3}$ | 1.87 | 2.41 | $\beta-\mathrm{CH}_{3}$ | 2.61 | 2.29 |
| $\mathrm{B} 11 \mathrm{H}_{3}$ | 2.25 | 2.41 | $\alpha-\mathrm{H}$ | 8.77 | 8.41 |
| B12H3 | 4.28 | 2.37 |  |  |  |

${ }^{a} \mathrm{CD}_{2} \mathrm{Cl}_{2}$. ${ }^{b}$ From parts of rings that are overlapped. Signals for $3,5-\mathrm{lut}^{\mathrm{t}}$ broad compared to those for 3,5 -luts.
one $\gamma-\mathrm{H}$, but two $\alpha-\mathrm{H}$ and two $\beta-\mathrm{CH}_{3}$ signals. For rapid $180^{\circ}$ rotation, 3,5 -lut would give one sharp signal for each proton type. Since the latter was found at room temperature (Figure 5), both $180^{\circ}$ rotation about the $\mathrm{Re}-\mathrm{N}$ bonds and the inversion of chirality are rapid.

On lowering the temperature, broadening is expected for the $\alpha-\mathrm{H}$ and $\beta-\mathrm{CH}_{3}$ signals but not for the $\gamma-\mathrm{H}$ signal, if only $\mathrm{Re}-\mathrm{N}$ rotation is slowed. However, if dynamic inversion of the chirality is slowed, broadening is expected for the $\gamma-\mathrm{H}$ signal as well as the other signals. Therefore, the temperature dependence of the spectrum of $\mathbf{2}$ was studied. The ${ }^{1} \mathrm{H}$ NMR spectrum of 2 had a similar, although less pronounced, temperature dependence as the spectrum for $\mathbf{1}$. A broadening of all peaks, including the $\gamma$-H signal, was observed below -40 ${ }^{\circ} \mathrm{C}$ (Figure 5). The broadening of the $\gamma$ - H signal is indicative of slowing but continued dynamic inversion of the chirality of 2, since $180^{\circ}$ rotation about the $\mathrm{Re}-\mathrm{N}$ bond will not influence the $\gamma$-H signal.

At still lower temperatures, e.g., $-90^{\circ} \mathrm{C}$, there are clearly two sets of 3,5-lut signals, and each has two $\alpha-\mathrm{H}$ and two $\beta-\mathrm{CH}_{3}$ signals. This result indicates that the $180^{\circ} \mathrm{Re}-\mathrm{N}$ rotation has also become slower. Also, the broadening of the two sets of 3,5-lut signals was not equal, indicating that one of the ligands was rotating more rapidly than the other. Thus 2 has two processes with rate-limiting steps of approximately the same time scale.

By tracing NOESY and EXSY cross-peaks in a 2D spectrum at $-100{ }^{\circ} \mathrm{C}$ (Figure 6), we can readily identify the signals for each $3,5-\mathrm{lut}$. For each $3,5-\mathrm{lut}$, one half has upfield and one half has downfield $\alpha-\mathrm{H}$ and $\beta-\mathrm{CH}_{3}$ signals. One 3,5-lut has all its respective signals upfield to those of the other. In other words, the more upfield of the $\alpha-\mathrm{H}$ signals of one is more upfield than the more upfield $\alpha-\mathrm{H}$ signal of the other. The assignment is consistent with the temperature dependence of the ${ }^{1} \mathrm{H}$ NMR spectrum of $\mathbf{2}$. This dependence allowed us to assign the signals to two sets because of the difference in broadening of each set; this difference is due to the differing rates of the $180^{\circ}$ rotation process for each 3,5 -lut type.

The structure of $\mathbf{2}$ should be similar to that of the analogous py dimer (3) ${ }^{14}$ since the 3,5 -lut methyl groups are pointing away from Re. The 3-type structure (see Figure 9 below) readily accounts for the shift pattern. As in the X-ray structure of $\mathbf{3}$, the plane of each $3,5-$ lut is expected to lie in a plane bisecting the cis bonds. As a result, half of each stacked 3,5-lut (3,5lut ${ }^{s}$ ) is stacked with and shielded by the corresponding half of the other 3,5-luts. (These halves are close to and related by the $C_{2}$-axis passing through the bridging O .) The signals which are most upfield are assigned to this half (Table 2). The other half of each $3,5-\mathrm{lut}^{\mathrm{s}}$ (displaced away from the $C_{2}$-axis) is closer to the terminal ligands, explaining the more downfield shift. From the py dimer structure, we expect one half of the terminal 3,5lut ( $3,5-\mathrm{lut}^{\text {t }}$ ) to lie over the ring of the Re-ligand-partner, $3,5-$ luts. This position leads to an upfield shift of the signals of
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Figure 9. Schematic representation similar to that in Figure 7, with L $=$ a symmetrical pyridine ligand. The rings of only pair 1 are shown to illustrate the relationships of the two halves of each symmetrical ligand going from a stacked site (left) to a terminal site (right).
this half of the $3,5-\mathrm{lut}$. Since this positioning is not so close as that between the two 3,5 -lut ${ }^{\mathrm{s}}$ 's, the upfield shift is smaller. With the assignments made, we can state that $180^{\circ}$ rotation is slightly faster for $3,5-\mathrm{lut}{ }^{\mathrm{t}}$ than for $3,5-\mathrm{lut}^{\mathrm{s}}$, since the signals of the former are broader at lower temperatures.

The presence of all possible EXSY peaks for $\gamma-\mathrm{H}, \alpha-\mathrm{H}$, and $\beta-\mathrm{CH}_{3}$ signals in the 2D NOESY/EXSY spectrum at $-100{ }^{\circ} \mathrm{C}$ (Figure 6) indicates that, even at this low temperature, inversion of chirality and $180^{\circ}$ rotation about the $\mathrm{Re}-\mathrm{N}$ bonds both occur to some extent. The exchange-NOEs in Figure 6 can be grouped into three classes, according to the type of information revealed. One class even reveals which halves of the Re-ligand-partners interchange (Figure 9).

Class (i) exchange-NOEs must result from inversion and reveal little about $\mathrm{Re}-\mathrm{N}$ rotation. Each $\gamma$ - H has cross-peaks to all four $\beta-\mathrm{CH}_{3}$ signals. Two are direct NOEs between the $\gamma$-H and the two $\beta-\mathrm{CH}_{3}$. The other two are class (i) exchangeNOEs, arising from the two direct NOEs after the ligand interchanges between t and s sites; these two are expected whether or not $180^{\circ} \mathrm{Re}-\mathrm{N}$ bond rotation is slow or fast.

The $\alpha-\mathrm{H}$ signals show direct NOEs to $\beta-\mathrm{CH}_{3}$ signals. At low contours, all $\alpha-\mathrm{H}$ signals also have two classes of exchangeNOEs. Class (ii) exchange-NOEs indicate only that $180^{\circ} \mathrm{Re}-\mathrm{N}$ rotation is occurring. For example, the downfield $\alpha-\mathrm{H}$ of $3,5-$ lut ${ }^{s}$ has, in addition to the direct NOE to the signal of the nearby $\beta-\mathrm{CH}_{3}$ (downfield 3,5-lut ${ }^{\mathrm{s}} \beta-\mathrm{CH}_{3}$ signal), an exchange-NOE to the upfield $\beta-\mathrm{CH}_{3}$ signal of $3,5-\mathrm{lut}$ s. This class (ii) exchangeNOE results from $\mathrm{Re}-\mathrm{N}$ rotation of $3,5-\mathrm{lut}$. The half of the 3,5-lut ${ }^{s}$ farther from the $C_{2}$-axis is placed closer to the axis by the rotation. Class (iii) exchange-NOEs can be explained only by the inversion of chirality without significant $180^{\circ} \mathrm{Re}-\mathrm{N}$ bond rotation. Examples of this class of exchange-NOE are the crosspeak from the upfield $\beta-\mathrm{CH}_{3}$ of $3,5-\mathrm{lut}{ }^{\mathrm{s}}$ to the downfield $\alpha-\mathrm{H}$ signal of the $3,5-l \mathrm{lut}$ and the cross-peak from the upfield $\alpha-\mathrm{H}$ of $3,5-\mathrm{lut}^{\mathrm{s}}$ to the downfield $\beta-\mathrm{CH}_{3}$ of $3,5-\mathrm{lut}$. These crosspeaks demonstrate that the Re-ligand-partners move in concert with the downfield half of the 3,5-lut ${ }^{t}$ undergoing physical exchange of position with the upfield half of the $3,5-l u t{ }^{\text {s }}$. This relationship establishes the exchange process shown in Figure 9 because inversion of chirality interchanges the positions of these halves of the two ligands. Thus, these class (iii) exchangeNOEs demonstrate that the overall net process, rotation about the $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ bond system and the $90^{\circ}$ rotation, is faster than $180^{\circ}$ rotation about the $\operatorname{Re}-\mathrm{N}(3,5-\mathrm{lut})$ bonds.

To aid the discussion, we define the coordination plane (CP); both CP's considered here will contain the ligand's N donor and the trans Cl atoms and are thus distinguished by the two other ligating atoms in the CP , namely $\mathrm{O}=\mathrm{Re}-\mathrm{O}$ (bridging) and (cis-N) $-\mathrm{Re}-\mathrm{Cl}$. Net inversion of chirality for $\left[\mathrm{Re}_{2} \mathrm{O}_{3} \mathrm{Cl}_{4} \mathrm{~L}_{4}\right]$ dimers with $C_{2}$-symmetrical ligands could occur if the ligand plane rotated by $\sim 90^{\circ}$ through either one of these two CP's (Figure 10). The sense of the rotation is different for the two options. The first option, rotation through the $\mathrm{O}=\mathrm{Re}-\mathrm{OCP}$


Figure 10. Schematic representation similar to that in Figure 8, but with a symmetrical pyridine-type ligand represented by the line. It is shown that the same form can be obtained by rotation through either CP.
(left side, Figure 10), explains the preferential interchange of the upfield and downfield halves of the 3,5-lut ligands (Figure 9 ) established by the class (iii) exchange-NOEs for 2 . The second option, rotation through the $\mathrm{Cl}-\mathrm{Re}-\mathrm{N} \mathrm{CP}$ (right side, Figure 10), is shown to be less favored by the exchange-NOE data.

The favored sense of rotation reflects the steric barriers L is expected to encounter as it rotates. The Cl groups and the N ligands are much more bulky than the O ligands (the Re complex at the other side of the bridge oxygen is too distant to affect this rotation in the anti conformer, see below). The barrier to rotation across this $\mathrm{Cl}-\mathrm{Re}-\mathrm{N} \mathrm{CP}$ would be higher. Furthermore, this second option does not hold for lopsided ligands, such as $\mathrm{Me}_{3} \mathrm{Bzm}$, since a different conformer would result if the plane of the ligand rotated past only the CP defined by the $\mathrm{Cl}-\mathrm{Re}-\mathrm{N}$ bonds (Figure 8). In order for $\mathbf{1}$ to undergo inversion of chirality by this route, each $\mathrm{Me}_{3} \mathrm{Bzm}$ would have to cross a $\mathrm{Cl}-\mathrm{Re}-\mathrm{NCP}$, a $\mathrm{O}=\mathrm{Re}-\mathrm{OCP}$, and then back across the $\mathrm{Cl}-$ $\mathrm{Re}-\mathrm{N} \mathrm{CP}$. Following this $270^{\circ}$ route to achieve a net $90^{\circ}$ $\mathrm{Re}-\mathrm{N}$ rotation (Figure 8) would greatly hinder the inversion rate; the bulky six-membered ring would have to clash severely with the cis $\mathrm{Me}_{3} \mathrm{Bzm}$ during the last $90^{\circ} \mathrm{Re}-\mathrm{N}$ rotation. Models suggest relatively much less steric hindrance for $90^{\circ}$ rotation of the L plane through the $\mathrm{O}=\mathrm{Re}-\mathrm{OCP}$, especially since the $\mathrm{Me}_{3} \mathrm{Bzm}$ five-membered rings are too small to clash.

The most likely route for the inversion of chirality that is consistent with the data for both $\mathbf{1}$ and $\mathbf{2}$ and that appears to be sterically feasible in the models involves three distinct types of rotation processes and an unstable intermediate. Thus, in the least hindered inversion pathway (Figure 11), the first process, a net $\sim 90^{\circ} \mathrm{Re}-\mathrm{O}-\mathrm{Re}$ rotation during which the stacked ligands are moved out and away from each other, forms an anti conformer with well-separated L's. Next, all four ligands rotate by $\sim 90^{\circ}$ in rapid succession past the respective $\mathrm{O}=\mathrm{Re}-\mathrm{OCP}$ 's (Figure 8). Finally there is a second net $\sim 90^{\circ} \mathrm{Re}-\mathrm{O}-\mathrm{Re}$ rotation in the same direction as the first $\mathrm{Re}-\mathrm{O}-\mathrm{Re}$ rotation.

The rates of inversion for $\mathbf{1}$ and $\mathbf{2}$ are expected to be similar for the route just described since there should be little difference in hindrance. However, as mentioned above, the pathway involving $270^{\circ}$ rotation should proceed much more slowly for 1. Thus, the rate of inversion is a way to assess the two pathways. Two distinct sets of $L$ signals were apparent by -60 ${ }^{\circ} \mathrm{C}$ for $\mathbf{1}$ (Figure 3) and $-80^{\circ} \mathrm{C}$ for 2 (Figure 5), but the shift differences are greater for $\mathbf{1}$. A comparison of the coalescence of signals with approximately the same difference in shifts for $L^{s}$ and $L^{t}$ suggests that the rate of inversion of chirality for $\mathbf{1}$ is less than an order of magnitude slower than for $\mathbf{2}$. This result, the exchange-NOE data, and examination of models all strongly suggest that the inversion occurs by the less hindered pathway with the L planes rotating by $90^{\circ}$ through the $\mathrm{O}=\mathrm{Re}-\mathrm{OCP}$ for both 1 and 2 (Figures 8 and 10).

## Conclusions

In octahedral sites in biological systems, lopsided ligands can select among four staggered orientations. Rarely has more than one staggered conformer been identified for any site. In model



Figure 11. Schematic representation of the complete mechanism proposed for the chirality inversion for $\mathbf{2}$. The chirality inversion occurs between the two chiral intermediates with the anti-conformation; these are shown at the bottom of the figure.
chiral octahedral complexes, a symmetrical ligand can select between only two possible staggered positions; the ligand can move between the two orientations by two opposing routes of rotation. In this study, we have shown how 2D exchange-NOEs can be used to distinguish these two routes of rotation. To our knowledge, this is the first time preferred pathways of rotation have been demonstrated. Exchange-NOEs have not been widely used to assess dynamic pathways. By comparisons with the analogous complex containing a lopsided ligand, we showed that the lopsided ligand very probably rotates using the same pathway as the symmetrical ligand. Such rotational motion can be significant in many biological processes such as in the formation of DNA cross-links by metal anticancer drugs.

The characteristic HH shift pattern and the NOE data demonstrate that there is only one significant conformer of $\mathbf{1}$ in solution and that this conformer has the HH,HT,HH structure found in the solid. This is the first case in which the predominant form of a complex with this type of untethered ligand has been shown to be HH in solution. The predominance of this form strongly suggests that the electrostatic attraction of the $\delta^{+}$aromatic proton in the imidazole ring of $\mathrm{Me}_{3} \mathrm{Bzm}^{\mathrm{t}}$ toward the negative central core of the dimer is responsible for these properties. This conclusion is supported by our studies with $\mathrm{Ru}^{\text {II }}$ complexes and other Re dimers. ${ }^{11,12,15}$

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Supporting Information Available: Crystallographic data for 1, including tables of positional parameters and bond distances and angles, anisotropic displacement coefficients, H -atom coordinates, view of the packing in $C 2 / c$, and a perspective view of the intermolecular stacking interactions in 1 (19 pages).

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